

**NOTE 1:** Answer all the 7 questions. Use the same symbols/notation given in the questions. Choose only one correct answer for questions 1 to 3.

**NOTE 2:** If you are sketching a graph, label the axes. No marks if axes are not labelled. No partial marks for questions 1 to 3. Write clearly and legibly.

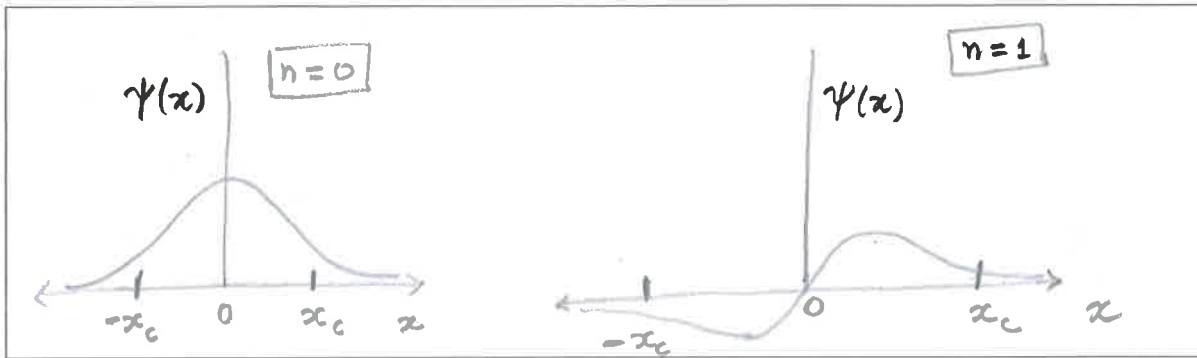
- For a parity-symmetric Hamiltonian in one dimension, which of the following statements is correct about degeneracy of bound states? (2)
  - Bound states are non-degenerate only for even states, odd states can be degenerate.
  - Bound states are always at least doubly degenerate due to parity.
  - Bound states are non-degenerate; each energy eigenvalue corresponds to an even or odd state.
  - Each energy eigenvalue corresponds to both an even and an odd eigenstate with the same energy.
  - Parity symmetry guarantees infinite degeneracy for each energy level.
- For the 1D harmonic oscillator, how does the Bohr correspondence principle manifest in the behavior of high- $n$  energy eigenstates  $\psi_n(x)$ , ( $n$  is the quantum number) ? (2)
  - The probability density  $|\psi_n(x)|^2$  approaches the classical probability distribution  $F(x)$ , concentrating in regions where the classical particle spends more time.
  - The high- $n$  eigenstates become sharply concentrated near  $x = 0$ , behaving like point particles around  $x = 0$ .
  - The spacing between adjacent quantised energy levels  $E_{n+1} - E_n$  increases without bound, matching the unbounded kinetic energy of a classical harmonic oscillator.
  - The probability density becomes completely uniform over all space, reflecting that the oscillator can be anywhere with equal likelihood.
  - Classical probability distribution  $F(x)$  is not related to  $|\psi_n(x)|^2$  since classical and quantum dynamics show very different features.
- Which of the statements is correct about time-reversal and time translation invariance ? (2)
  - Time translation invariance of  $H$  requires that  $[H, L_z] = 0$ .
  - If  $\psi(x, t)$  is a complex quantum state and it satisfies time-dependent Schrodinger equation, then  $\psi(x, -t)$  will also be satisfied by the same Schrodinger equation and hence the system is time-reversal invariant.
  - Time-reversal invariance requires a Hamiltonian  $H$  such that  $H^* = H$ .
  - Time-reversal invariance is the property of eigenstates of the system.
  - Time translation invariance means that the initial states do not evolve when translated from  $t$  to  $t + \delta t$ .

For the MCQs 1-3 given above, write your answer only inside the three boxes here:

1       2       3

4. Consider the oscillator  $H = p^2/2m + m\omega^2 x^2/2$ .

(a) Sketch the  $n = 0$  and  $n = 1$  state of the harmonic oscillator ( $n$  is the quantum number). On the same graph, show the location of the classical turning points  $x_c$  for both the states. (3)



(b) Write (do not derive) an expression for the classical turning point  $x_c$  in terms of  $n$  and  $\hbar$ . (2)

$$x_c = \pm \sqrt{(2n+1) \frac{\hbar}{m\omega}}$$

5. Compute  $\langle m | \hat{X} | n \rangle$ , where  $m$  and  $n$  are harmonic oscillator states. Take  $\hbar/2m\omega = 1$ . (3)

6. Let  $H = p^2/2m + \beta$ , where  $\beta > 0$  is some constant, be a system that is invariant under translation in 1D. If  $\hat{U} = e^{i\hat{H}t/\hbar}$  is used to time evolve some state, then show that  $\hat{U}$  is also invariant under translation. (3)

7. Let  $\hat{U}$  be the operator for infinitesimal rotation by an angle  $\epsilon$  about  $z$ -axis, and  $\hat{X}$  is the position operator. Then, starting from  $\hat{U}^\dagger \hat{X} \hat{U} = \hat{X} - \epsilon \hat{Y}$ , determine the commutator  $[\hat{X}, \hat{L}_z]$ . (3)

Start writing your answers for questions 5 to 7 below

5. position operator:  $\hat{X} = \sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger)$

$$a|n\rangle = \sqrt{n}|n-1\rangle \quad \text{and} \quad a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle.$$

$$\therefore \langle m | \hat{X} | n \rangle = \sqrt{\frac{\hbar}{2m\omega}} (\langle m | a | n \rangle + \langle m | a^\dagger | n \rangle)$$

$$= \sqrt{\frac{\hbar}{2m\omega}} (\sqrt{n} \langle m | n-1 \rangle + \sqrt{n+1} \langle m | n+1 \rangle)$$

$$= \sqrt{\frac{\hbar}{2m\omega}} (\sqrt{n} \delta_{m, n-1} + \sqrt{n+1} \delta_{m, n+1})$$

$$= (\sqrt{n} \delta_{m, n-1} + \sqrt{n+1} \delta_{m, n+1}) \quad \text{Since } \frac{\hbar}{2m\omega} = 1.$$

$$6. \quad H = \frac{p^2}{2m} + \beta \quad (\beta > 0).$$

$T$  is the translation operator such that  $\hat{T} |x\rangle = |x+\epsilon\rangle$

If  $H$  is invariant under spatial translation, then  $[H, \hat{T}] = 0$

If  $\hat{U} = e^{-iHt/\hbar}$ , then we need to show that  $[T, \hat{U}] = 0$

To show this, write  $\hat{U} = \sum_{n=0}^{\infty} \frac{(-it/\hbar)^n}{n!} H^n$

We need to show that  $[T, U] = 0$

$$\text{Now, } [\hat{U}, \hat{T}] = \sum_{n=0}^{\infty} \frac{(-it/\hbar)^n}{n!} [H^n, T]$$

Implies that we must show that  $[H^n, T] = 0$  for all  $n$ , given that it is true for  $n=1$ .

By induction: Assume  $[H^k, T] = 0$  for  $k \leq n$ .

$$\begin{aligned} \text{Then, } [H^{n+1}, T] &= H^n [H, T] + [H^n, T] H \\ &= H^n \cdot 0 + 0 \cdot H = 0 \end{aligned}$$

$\therefore [H^n, T] = 0$  for all  $n$ . Hence  $[\hat{U}, \hat{T}] = 0$ .

Another method: we need to show  $[U, T] = e^{-iHt/\hbar} T - T e^{-iHt/\hbar} = 0$

$$\text{we } e^A B = \sum_n \frac{A^n}{n!} B. \quad \text{If } [A, B] = 0, \text{ then}$$

$$e^A B = B e^A.$$

Hence,  $[U, T] = 0$ .

$$7. \quad \hat{U} = I - \frac{i\varepsilon}{\hbar} L_z$$

$$U^\dagger X U = X - \varepsilon Y$$

$$\left( I + \frac{i\varepsilon}{\hbar} L_z \right) X \left( I - \frac{i\varepsilon}{\hbar} L_z \right) = X - \varepsilon Y$$

$$X - \frac{i\varepsilon}{\hbar} X L_z + \frac{i\varepsilon}{\hbar} L_z X = X - \varepsilon Y$$

$$\frac{i\varepsilon}{\hbar} (L_z X - X L_z) = -\varepsilon Y$$

$$\frac{i\varepsilon}{\hbar} (X L_z - L_z X) = \varepsilon Y$$

$$\therefore [X, L_z] = \frac{\hbar}{i} Y \quad \text{or} \quad [X, L_z] = -i\hbar Y$$

$$\text{or} \quad [L_z, X] = i\hbar Y$$

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