

x Periodic perturbation

example: atom interacting with monochromatic light beam
atom placed between plates of a capacitor
connected to ac source

Let $H'(t) = H' e^{-i\omega t}$

Our system begins to be perturbed by H' at time $t=0$.

Transition amplitude from $|i^0\rangle$ to $|f^0\rangle$ in time t is

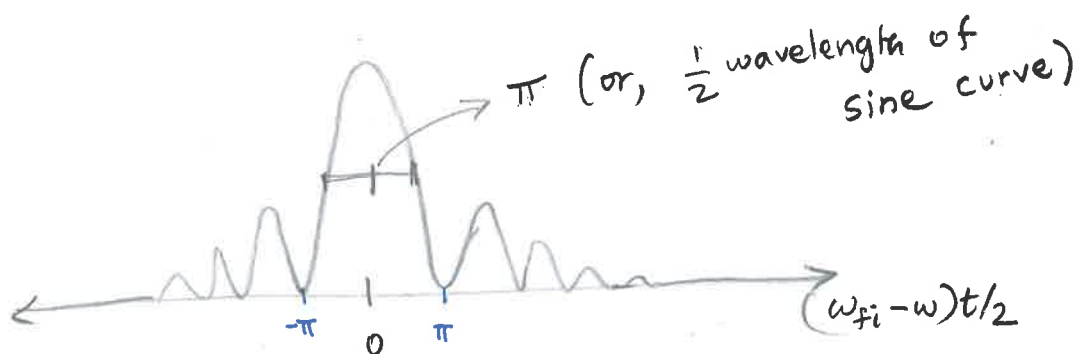
$$d_f(t) = \left(-\frac{i}{\hbar}\right) \int_0^t \langle f^0 | H' | i^0 \rangle e^{-i\omega t'} e^{i\omega_{fi} t'} dt'$$

$$= -\frac{i}{\hbar} \langle f^0 | H' | i^0 \rangle \left. \frac{e^{i(\omega_{fi} - \omega)t'}}{i(\omega_{fi} - \omega)} \right|_0^t$$

$$= -\frac{1}{\hbar} \langle f^0 | H' | i^0 \rangle \frac{[e^{i(\omega_{fi} - \omega)t} - 1]}{(\omega_{fi} - \omega)}$$

$$d_f(t) = -\frac{1}{\hbar} \langle f^0 | H' | i^0 \rangle e^{i(\omega_{fi} - \omega)t/2} \left[\frac{e^{i(\omega_{fi} - \omega)t/2} - e^{-i(\omega_{fi} - \omega)t/2}}{(\omega_{fi} - \omega)t/2} \right] \frac{t}{2}$$

$$|d_f(t)|^2 = \frac{1}{\hbar^2} |\langle f^0 | H' | i^0 \rangle|^2 \left(\frac{\sin(\omega_{fi} - \omega)t/2}{(\omega_{fi} - \omega)t/2} \right)^2 t^2 \quad \text{--- (1)}$$



$$\therefore \left| (\omega_{fi} - \omega) \frac{t}{2} \right| \lesssim \pi$$

System will go to states $|f\rangle$ under this condition

$$\left(\frac{E_f^0 - E_i^0}{\hbar} \right) \frac{t}{2} = \pm \pi + \frac{\omega t}{2}$$

$$E_f^0 t - E_i^0 t = \pm 2\pi\hbar + \hbar\omega t$$

$$\therefore E_f^0 t = (E_i^0 t + \hbar\omega t) \pm 2\pi\hbar$$

$$\text{or, } E_f^0 - E_i^0 = \hbar\omega \left(1 \pm \frac{2\pi}{\omega t} \right)$$

- For small t , E_f^0 is not the preferred final state.
- For $\omega t \gg 2\pi$, $E_f^0 = E_i^0 + \hbar\omega$ is the most favoured.
We need few cycles of perturbation for this to happen.

What happens as $t \rightarrow \infty$?

Consider system exposed to perturbation from $t = -\frac{T}{2}$ to $\frac{T}{2}$ and $T \rightarrow \infty$. Then,

$$d_f = \lim_{T \rightarrow \infty} \frac{2t}{\hbar} \int_{-T/2}^{T/2} H'_{fi} e^{i(\omega_{fi} - \omega)t'} dt'$$

$$= -\frac{i}{\hbar} H'_{fi} 2\pi \delta(\omega_{fi} - \omega)$$

$$P_{i \rightarrow f} = |d_f|^2 = \frac{4\pi^2}{\hbar^2} |H'_{fi}|^2 \delta(\omega_{fi} - \omega) \delta(\omega_{fi} - \omega)$$

$$\delta(\omega_{fi} - \omega) \delta(\omega_{fi} - \omega) = \lim_{T \rightarrow \infty} \delta(\omega_{fi} - \omega) \frac{1}{2\pi} \int_{-T/2}^{T/2} e^{i(\omega_{fi} - \omega)t} dt$$

$$= \lim_{T \rightarrow \infty} \delta(\omega_{fi} - \omega) \frac{T}{2\pi}$$

$$\therefore P_{i \rightarrow f} = \frac{4\pi^2}{\hbar^2} |H'_{fi}|^2 \frac{T}{2\pi} \delta(E_f^0 - E_i^0 - \hbar\omega) \cdot \hbar$$

Transition
rate

$$R_{i \rightarrow f} = \frac{P_{i \rightarrow f}}{T} = \frac{2\pi}{\hbar} |\langle f^0 | H' | i^0 \rangle|^2 \delta(E_f^0 - E_i^0 - \hbar\omega)$$

This is Fermi's Golden rule. Has many applications.

- Due to δ -function in the formula, transition rate makes sense only under an integral.
- As $t \rightarrow \infty$, Fourier transform of the perturbation is sharply at the frequency of perturbation (due to δ -function). System "sees" a single frequency.

In short-time limit: use Eq (1), and expand sine term

$$P_{i \rightarrow f} = \frac{t^2 |H'_{fi}|^2}{\hbar^2}$$

$$\text{Transition rate: } R_{i \rightarrow f} = \frac{t}{\hbar^2} |H'_{fi}|^2$$

Excited states lie in a band

$$dN = g(E_k) dE_k$$

$\rightarrow$ density of final states

The prob. of transition occurs to a state in a band of width 2Δ centred at E_k

$$\bar{P}_{i \rightarrow f} = \int_{E_f - \Delta}^{E_f + \Delta} P_{i \rightarrow f} g(E_f) dE_f$$

Substitute for $P_{i \rightarrow f}$ from Eq (1):

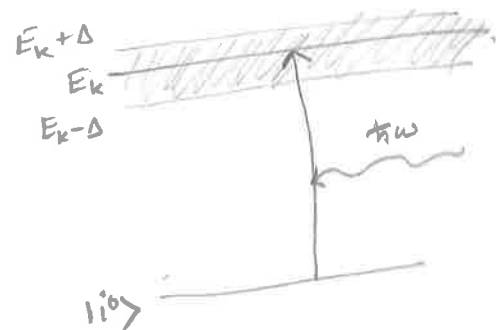
$$\bar{P}_{i \rightarrow f} = \int_{E_f - \Delta}^{E_f + \Delta} dE_f g(E_f) \frac{|\langle f^0 | H' | i^0 \rangle|^2}{\hbar^2} \frac{\sin^2 \beta}{\beta^2/t^2}$$

$$\text{where } 2\hbar\beta = \hbar(\omega_{fi} - \omega)t = (E_f - E_i - \hbar\omega)t$$

For fixed $|i^0\rangle$, t and ω , we have $dE_f = \frac{2\hbar}{t} d\beta$

$$\bar{P}_{i \rightarrow f} = \frac{2t}{\hbar} \int_{-\delta}^{\delta} g(E_f) |H'_{fi}|^2 \frac{\sin^2 \beta}{\beta^2} d\beta,$$

δ is the value of β corresponding to Δ for E_f .



$$\bar{P}_{i \rightarrow f} = \frac{2t}{\hbar} g(E_f) |H'_{fi}|^2 \int_{-\infty}^{\infty} \frac{\sin^2 \beta}{\beta^2} d\beta$$

assumption: $g(E_f)$ and $|H'_{fi}|^2$ vary slowly with β .

$$= \left[\frac{2\pi}{\hbar} g(E_f) |H'_{fi}|^2 \right] t$$

The associated transition probability rate.

$$\bar{R}_{i \rightarrow f} = \frac{2\pi}{\hbar} g(E_f) |H'_{fi}|^2$$