

- 1. $\hat{P}^2 \Psi(x) = \Psi(x)$. Hence, eigenvalues are ± 1 .

$$\begin{aligned}\hat{P}(e^{-x^2/2} \cos x) &= e^{-(-x)^2/2} \cos(-x) \\ &= e^{-x^2/2} \cos x\end{aligned}$$

\therefore Eigenvalue is +1. Hence, even parity.

- 2. $E = 2k + |m| + 1 , \quad k = 0, 1, 2, \dots$

$$\text{If } n = 2k + |m|$$

For a given n , $k = 0, 1, 2, \dots n-1$.

Degeneracy is $n+1$.

- 3. Let $H \Psi(x) = E \Psi(x)$. Note $\hat{P} \Psi(x) = \Psi(-x)$

$$\hat{P} H \Psi(x) = E \hat{P} \Psi(x)$$

Since H and \hat{P} commute, we have

$$H \hat{P} \Psi(x) = E \hat{P} \Psi(x)$$

$$\therefore H \Psi(-x) = E \Psi(-x)$$

Thus, $\Psi(-x)$ is also an eigenstate with same eigenvalue.

■ 4. $\delta(x^2 - a^2) = \frac{1}{2a} [\delta(x-a) + \delta(x+a)]$

$$\int_{-\infty}^{\infty} \delta(x^2 - a^2) f(x) dx = \int_{-\infty}^0 \delta(x^2 - a^2) f(x) dx + \int_0^{\infty} \delta(x^2 - a^2) f(x) dx \quad \dots \text{--- } ①$$

$f(x) \rightarrow$ some probe function

$$\left| \begin{array}{l} \text{put } x = -\sqrt{u}, \quad dx = -\frac{du}{2\sqrt{u}} \\ \therefore x^2 = u \end{array} \right.$$

Then,

$$\int_{-\infty}^0 \delta(x^2 - a^2) f(x) dx = -\frac{1}{2} \int_{\infty}^0 \delta(u - a^2) \frac{f(-\sqrt{u})}{\sqrt{u}} du$$

$$= \frac{1}{2} \int_0^{\infty} \delta(u - a^2) \frac{f(-\sqrt{u})}{2\sqrt{u}} du = \frac{1}{2a} f(-a)$$

We have used the identity $\int g(y) \delta(y - \alpha) dy = g(\alpha)$

Consider the 2nd term in Eqn (1) : $\int_0^{\infty} \delta(x^2 - a^2) f(x) dx$

$$\text{Put } x = \sqrt{u}, \quad dx = \frac{du}{2\sqrt{u}}$$

$$x^2 = u$$

$$\int_0^{\infty} \delta(x^2 - a^2) f(x) dx = \int_0^{\infty} \delta(u - a^2) \frac{f(\sqrt{u})}{2\sqrt{u}} du$$

$$= \frac{1}{2a} f(a)$$

$$\therefore \int_{-\infty}^{\infty} \delta(x^2 - a^2) f(x) dx = \frac{1}{2a} [f(-a) + f(a)]$$

note that $f(a) = \int f(x) \delta(x-a) dx$

$$\int_0^{\infty} \delta(x^2 - a^2) f(x) dx = \frac{1}{2a} \left[\int f(x) \delta(x+a) dx + \int f(x) \delta(x-a) dx \right]$$

In "unofficial" notation:

$$\delta(x^2 - a^2) = \frac{1}{2a} [\delta(x+a) + \delta(x-a)]$$

■ 5. $\hat{U}_{[R(\varepsilon)]}^\dagger P_y U_{[R(\varepsilon)]} = \left(I + \frac{i\varepsilon}{\hbar} L_z \right) P_y \left(I - \frac{i\varepsilon}{\hbar} L_z \right)$

$$= P_y + \frac{i\varepsilon}{\hbar} (L_z P_y - P_y L_z) = P_y + \frac{i\varepsilon}{\hbar} [L_z, P_y]$$

$$= P_y + \frac{i\varepsilon}{\hbar} [L_z, P_y]$$

$$= P_y - \frac{i\varepsilon}{\hbar} [P_y, L_z] \quad \text{----- (2)}$$

$$= P_y - \frac{i\varepsilon}{\hbar} [P_y (x P_y - y P_x) - (x P_y - y P_x) P_y]$$

$$= P_y - \frac{i\varepsilon}{\hbar} \left(P_y x P_y - P_y y P_x - x P_y P_y + y P_x P_y \right)$$

- Used the fact that $P_x P_y = P_y P_x$,
- and $[x, P_y] = 0$. Hence, the cancellation.

We are left with

$$\begin{aligned}
 &= P_y - \frac{i\varepsilon}{\hbar} (y P_x P_y - P_y y P_x) \\
 &= P_y - \frac{i\varepsilon}{\hbar} (y P_y - P_y y) P_x \\
 &= P_y - \frac{i\varepsilon}{\hbar} [y, P_y] P_x = P_y - \frac{i\varepsilon}{\hbar} (i\hbar) P_x \\
 &= P_y + \varepsilon P_x \quad \text{--- --- (3)}
 \end{aligned}$$

Comparing Eq(2) and Eq(3), we get

$$-\frac{i}{\hbar} [P_y, L_z] = P_x$$

This gives $[P_y, L_z] = i\hbar P_x$.