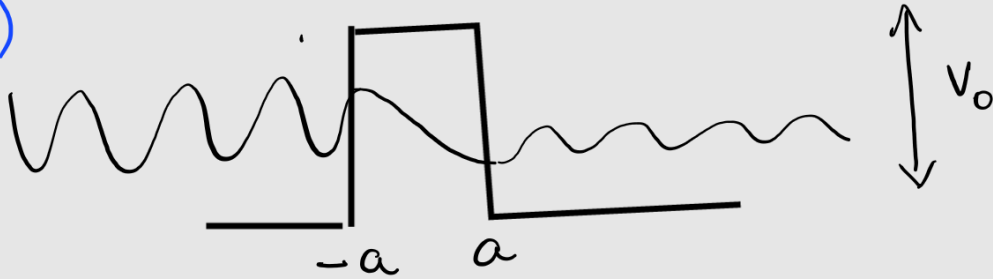


1a)



b) In the limit $V_0 \rightarrow \infty$, $R=1$ and $T=0$.

2a)

conserved quantities

Energy E Total angular momentum squared L^2 z-component of angular momentum L_z

Commutation relations:

$$[H, H] = 0, \quad [H, L^2] = 0, \quad [H, L_z] = 0.$$

b) $m = -l, -l+1, \dots, 0, 1, 2, \dots, l$

degree of degeneracy: $2l+1$

3)

$$\langle x \rangle = \int_0^\infty |A|^2 x e^{-|x|/\Delta} dx$$

$$\langle x \rangle = |A|^2 \Delta^2 \int_0^\infty \frac{x}{\Delta} e^{-x/\Delta} \frac{dx}{\Delta}$$

put $u = x/\Delta$

$$\langle x \rangle = |A|^2 \Delta^2 \underbrace{\int_0^\infty u e^{-u} du}_1$$

$$\therefore \langle x \rangle = |A|^2 \Delta^2$$

$$4) [L_x, L_y]_+ = L_x L_y + L_y L_x$$

$$\text{put } L_x = \frac{1}{2}(L_+ + L_-), \quad L_y = \frac{1}{2i}(L_+ - L_-)$$

Now,

$$L_x L_y + L_y L_x = \frac{1}{4i} \left\{ (L_+ + L_-)(L_+ - L_-) + (L_+ - L_-)(L_+ + L_-) \right\}$$

$$= \frac{1}{4i} \left\{ L_+^2 - L_-^2 + L_+^2 - L_-^2 \right\}$$

$$= \frac{1}{2i} (L_+^2 - L_-^2) = -\frac{i}{2} (L_+^2 - L_-^2)$$

$$5) A = \begin{pmatrix} 1 & 3 \\ 5 & 4 \end{pmatrix}$$

$$A = m_0 I + m_1 \sigma_1 + m_2 \sigma_2 + m_3 \sigma_3$$

$$m_i = \frac{1}{2} \text{Tr}(A \sigma_i)$$

Using this, we get $m_0 = \frac{5}{2}$, $m_1 = 4$, $m_2 = -i$, $m_3 = -\frac{3}{2}$

$$6) \quad \hat{p} = i \sqrt{\frac{m\omega\hbar}{2}} (a^\dagger - a)$$

$$p^2 = -\frac{m\omega\hbar}{2} (a^{\dagger 2} - a^\dagger a - a a^\dagger - a^2)$$

$$\langle 0 | p^2 | 0 \rangle = \frac{m\omega\hbar}{2} \langle 0 | a^2 + a^\dagger a + a a^\dagger - a^{\dagger 2} | 0 \rangle$$

We know: $a|n\rangle = \sqrt{n}|n-1\rangle$, $a^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$

$$\begin{aligned} \bullet \langle 0 | a^2 | 0 \rangle &= 0 & \bullet \langle 0 | a a^\dagger | 0 \rangle &= \langle 0 | a | 1 \rangle \\ & & &= \langle 0 | 0 \rangle = 1 \end{aligned}$$

$$\bullet \langle 0 | a^\dagger a | 0 \rangle = 0$$

$$\begin{aligned} \bullet \langle 0 | a^{\dagger 2} | 0 \rangle &= \langle 0 | a^\dagger | 1 \rangle \\ &= \langle 0 | 2 \rangle \sqrt{2} = 0 \end{aligned}$$

$$\therefore \langle 0 | p^2 | 0 \rangle = \frac{m\omega\hbar}{2}$$

7)

Given that $\psi_{n,n-1,m} = A_n r^{n-1} e^{-r/na_0} Y_{n-1}^m(\theta, \varphi)$

Probability to find electron
in r and $r+dr$ is

$$\begin{aligned} P(r) &= A_n^2 r^{2(n-1)} e^{-2r/na_0} r^2 \\ &= A_n^2 r^{2n} e^{-2r/na_0} \end{aligned}$$

r^2 is multiplied to account for Jacobian.

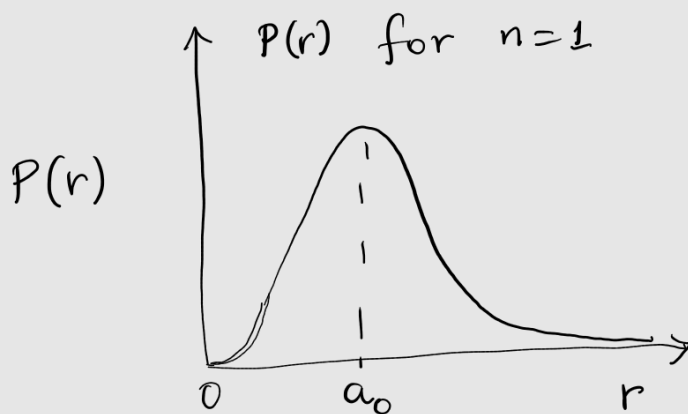
To find maximum of $P(r)$

$$\frac{d}{dr} P(r) = 0$$

$$\Rightarrow A_n \left[2n r^{2n-1} e^{-2r/na_0} + r^{2n} \left(-\frac{2}{na_0} \right) e^{-2r/na_0} \right] = 0$$

$$\frac{2n}{r} - \frac{2}{na_0} = 0 \Rightarrow r = n^2 a_0$$

$\therefore$ Probability to find electron is maximum at $n^2 a_0$



$$8) \quad \langle L_x \rangle = \langle l m | L_x | l m \rangle$$

$$L_x = \frac{L_+ + L_-}{2}$$

$$\langle L_x \rangle = \frac{1}{2} \langle l m | L_+ + L_- | l m \rangle$$

$$L_{\pm} |l, m\rangle = \hbar \sqrt{(l \mp m)(l \pm m + 1)} |l, m \pm 1\rangle$$

$$\langle L_x \rangle = \frac{1}{2} \hbar \sqrt{(l-m)(l+m+1)} \langle l, m | l, m+1 \rangle \\ + \frac{1}{2} \hbar \sqrt{(l+m)(l-m+1)} \langle l, m | l, m-1 \rangle$$

$$\langle l, m | l, m+1 \rangle = \langle l, m | l, m-1 \rangle = 0$$

$$\langle L_x \rangle = 0$$

$$\langle L_x^2 \rangle = \left\langle \left(\frac{L_+ + L_-}{2} \right)^2 \right\rangle = \frac{1}{4} \langle l, m | L_+^2 + L_+ L_- + L_- L_+ + L_-^2 | l, m \rangle$$

$$\langle l, m | L_+^2 | l, m \rangle = \langle l, m | L_-^2 | l, m \rangle = 0$$

$$\langle l, m | L_+ L_- | l, m \rangle = \hbar \sqrt{(l+m)(l-m+1)} \langle l, m | L_+ | l, m-1 \rangle \\ = \hbar \sqrt{(l+m)(l-m+1)} \hbar \sqrt{(l-m+1)(l+m)} \langle l, m | l, m \rangle$$

$$\langle L_+ L_- \rangle = \hbar^2 (l+m)(l-m+1)$$

$$\langle l, m | L_- L_+ | l, m \rangle = \hbar \sqrt{(l-m)(l+m+1)} \langle l, m | L_- | l, m+1 \rangle \\ = \hbar \sqrt{(l-m)(l+m+1)} \hbar \sqrt{(l+m+1)(l-m)} \langle l, m | l, m \rangle$$

$$\langle L_- L_+ \rangle = \hbar^2 (l-m)(l+m+1)$$

$$\langle L_x^2 \rangle = \frac{1}{4} \hbar^2 \left[(l+m)(l-m+1) + (l-m)(l+m+1) \right]$$

$$= \frac{\hbar^2}{4} 2(l(l+1) - m^2)$$

$$\langle L_x^2 \rangle = \frac{\hbar^2}{2} (l(l+1) - m^2)$$

$$\sigma_{L_x} = \sqrt{\langle (L_x - \langle L_x \rangle)^2 \rangle} = \sqrt{\langle L_x^2 \rangle}$$

$$\sigma_{L_x} = \frac{\hbar}{\sqrt{2}} \sqrt{l(l+1) - m^2}$$

9 a) Rotation operator for spin

$$U(\theta) = I \cos \frac{\theta}{2} - i \sin \frac{\theta}{2} \hat{\theta} \cdot \vec{\sigma}$$

$$\vec{\sigma} = \sigma_x \hat{i} + \sigma_y \hat{j} + \sigma_z \hat{k}$$

$$\hat{\theta} = \hat{i} \text{ (for this problem), } \theta = \pi/2$$

Then, we get

$$U\left(\frac{\pi}{2}\right) = I \cos \frac{\pi}{4} - i \sin \frac{\pi}{4} \sigma_x$$

$$= \frac{1}{\sqrt{2}} I - \frac{i}{\sqrt{2}} \sigma_x$$

Note $|\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$U\left(\frac{\pi}{2}\right) = \frac{1}{\sqrt{2}} \left[\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ -i & 1 \end{pmatrix}$$

we need :

$$\begin{aligned} U\left(\frac{\pi}{2}\right) |\uparrow\rangle &= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ -i & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ -i \end{pmatrix} \\ &= \begin{pmatrix} 1 \\ -i \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} - i \begin{pmatrix} 0 \\ 1 \end{pmatrix} \\ &= |\uparrow\rangle - i |\downarrow\rangle \end{aligned}$$

$$\blacksquare U\left(\frac{\pi}{2}\right) |\uparrow\rangle = |\uparrow\rangle - i |\downarrow\rangle$$

9b) Let $|u\rangle = \begin{pmatrix} a \\ b \end{pmatrix}$.

$$\begin{aligned} \langle \sigma_x \rangle &= \langle u | \sigma_x | u \rangle = \begin{pmatrix} a^* & b^* \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} \\ &= (a^* - b^*) \begin{pmatrix} a \\ b \end{pmatrix} = |a|^2 - |b|^2 \end{aligned}$$

$$\langle \sigma_x \rangle = 0 \text{ implies that } |a|^2 - |b|^2 = 0$$

$$\text{Hence, we can choose } a = e^{i\theta_1}, \quad b = e^{i\theta_2}$$

$$\therefore |u\rangle = \begin{pmatrix} e^{i\theta_1} \\ e^{i\theta_2} \end{pmatrix}$$

$$\begin{aligned} \langle \sigma_y \rangle &= \langle u | \sigma_y | u \rangle = \begin{pmatrix} e^{-i\theta_1} & e^{-i\theta_2} \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} e^{i\theta_1} \\ e^{i\theta_2} \end{pmatrix} \\ &= \begin{pmatrix} i e^{-i\theta_2} & -i e^{-i\theta_1} \end{pmatrix} \begin{pmatrix} e^{i\theta_1} \\ e^{i\theta_2} \end{pmatrix} \end{aligned}$$

$$= i \left[e^{i(\theta_1 - \theta_2)} - e^{-i(\theta_1 - \theta_2)} \right]$$

$$\langle \sigma_y \rangle = -2 \sin(\theta_1 - \theta_2)$$

We require that $\langle \sigma_y \rangle = 0$ or $\sin(\theta_1 - \theta_2) = 0$

$$\text{Now, } \langle \sigma_x \rangle = \begin{pmatrix} e^{-i\theta_1} & e^{-i\theta_2} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{i\theta_1} \\ e^{i\theta_2} \end{pmatrix}$$

$$\langle \sigma_x \rangle = \begin{pmatrix} e^{-i\theta_2} & e^{-i\theta_1} \end{pmatrix} \begin{pmatrix} e^{i\theta_1} \\ e^{i\theta_2} \end{pmatrix}$$

$$= 2 \cos(\theta_1 - \theta_2)$$

We require that $\langle \sigma_x \rangle = 0$ or $\cos(\theta_1 - \theta_2) = 0$

Now, if $\theta_1 - \theta_2 = \beta$, then both conditions

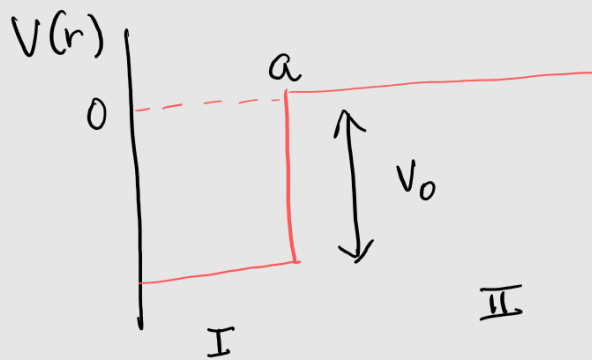
$$\langle \sigma_y \rangle = \sin \beta = 0 \quad \text{and}$$

$$\langle \sigma_x \rangle = \cos \beta = 0 \quad \text{CANNOT be simultaneously satisfied.}$$

Hence, it is not possible to simultaneously satisfy

$$\langle \sigma_x \rangle = \langle \sigma_y \rangle = \langle \sigma_z \rangle = 0.$$

10) Potential



- Radial Schrodinger equation

$$-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \left[V(r) + \frac{\hbar^2}{2m} \frac{l(l+1)}{r^2} \right] u = E u$$

- We require $l=0$ case.

$$\therefore -\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + V(r) u = E u$$

- Applied to this problem, we have

$$-\frac{\hbar^2}{2m} \frac{d^2 u_I}{dr^2} - V_0 u_I = E u_I \quad \text{for } (r \leq a) \quad \text{region I}$$

$$-\frac{\hbar^2}{2m} \frac{d^2 u_{II}}{dr^2} = E u_{II} \quad \text{for } (r > a) \quad \text{region II}$$

- Consider region I. Here $E < 0$. $r \leq a$

$$\frac{d^2 u_I}{dr^2} + \frac{2m V_0}{\hbar^2} u_I = -\frac{2m}{\hbar^2} (-E) u_I$$

$$\frac{d^2 u_I}{dr^2} + \frac{2m}{\hbar^2} (V_0 - E) u_I = 0$$

$$\text{put } k_1^2 = \frac{2m}{\hbar^2} (V_0 - E)$$

$$\therefore \frac{d^2 u_I}{dr^2} + k_1^2 u_I = 0 \quad \text{--- --- --- (1)}$$

■ Consider region II. Here too $E < 0$. $r > a$

$$-\frac{\hbar^2}{2m} \frac{d^2 u_{II}}{dr^2} = (-E) u_{II}$$

$$\Rightarrow \frac{d^2 u_{II}}{dr^2} - \frac{2mE}{\hbar^2} u_{II} = 0 \Rightarrow \frac{d^2 u_{II}}{dr^2} - k_2^2 u_{II} = 0 \quad \text{--- --- (2)}$$

$$k_2^2 = 2mE/\hbar^2$$

■ Solution in region I ($r \leq a$)

$$u_I(r) = A \sin k_1 r + B \cos k_1 r$$

Since we need $u_I(r) = 0$ at $r = 0$

■ Solution in region II ($r > a$)

$$u_{II}(r) = C e^{k_2 r} + D e^{-k_2 r}$$

Match solutions at $r = a$.

$$u_I(a) = u_{II}(a)$$

$$A \sin k_1 a = D e^{-k_2 a}$$

Match derivatives at $r = a$:

$$A k_1 \cos k_1 a = -D k_2 e^{-k_2 a}$$

$$\text{put } \frac{A}{D} = C$$

$\therefore$ we have

$$C \sin k_1 a = e^{-k_2 a} \quad \text{and} \quad C k_1 \cos k_1 a = -k_2 e^{-k_2 a}$$

Dividing the two, we get quantisation condition:

$$\boxed{\tan k_1 a = -\frac{k_1}{k_2}} \quad \text{--- (3)}$$

b) There will be no energy levels if condition in Eq (3) is violated.

i.e if $\tan k_1 a > 0$

That is, $k_1 a < \pi/2$ or $k_1^2 a^2 < \frac{\pi^2}{4}$

$$\text{i.e } \frac{2m}{\hbar^2} (V_0 - E) a^2 < \frac{\pi^2}{4}$$

$$\frac{2m V_0 a^2}{\hbar^2} < \frac{\pi^2}{4} + \frac{2m E a^2}{\hbar^2}$$

$$\Rightarrow V_0 a^2 < \frac{\pi^2 \hbar^2}{8m}$$